Scaling in a network of Zo strings

Mark Hindmarsh (Sussex, Helsinki), Kari Rummukainen (Helsinki) and David J. Weir (Stavanger)

What was known

- Spontaneous breaking of a non-abelian symmetry can produce 'beads' consisting of 't Hooft-Polyakov monopoles, on cosmic strings.
- It is not known how the monopoles influence the dynamics of the resulting string network.

What this work adds

- We have carried out simulations of this scenario for the first time.
- Monopoles are carried along by the strings; the network behaves similarly to an abelian string network.

Next steps

- Study larger ratios between the monopole and string scales, to check whether monopoles eventually slow strings down.
- Observational predictions for strings in grand unified models, e.g. SO(10).

Introduction

▶ We study the formation of cosmic string networks in the model with Lagrangian ^{1,2}

$$\mathcal{L} = -\frac{1}{4} F^{a}_{\mu\nu} F^{\mu\nu a} + \sum_{n} \operatorname{Tr}[D_{\mu}, \Phi_{n}] [D^{\mu}, \Phi_{n}] - V(\Phi_{1}, \Phi_{2}); \qquad V(\Phi_{1}, \Phi_{2}) = -m_{1}^{2} \operatorname{Tr} \Phi_{1}^{2} - m_{2}^{2} \operatorname{Tr} \Phi_{2}^{2} + \lambda \left[\left(\operatorname{Tr} \Phi_{1}^{2} \right)^{2} + \left(\operatorname{Tr} \Phi_{2}^{2} \right)^{2} \right] + \kappa \left(\operatorname{Tr} \Phi_{1} \Phi_{2} \right)^{2}$$

where $D_{\mu} = \partial_{\mu} + igA_{\mu}$, $F_{\mu\nu} = F^{a}_{\mu\nu}\tau^{a}$ and $A_{\mu} = A^{a}_{\mu}\tau^{a}$, $\tau^{a} = \sigma^{a}/2$. Here, Φ_{1} and Φ_{2} are adjoint Higgs fields ($\Phi = \phi^{a}\sigma^{a}$).

• The system undergoes two symmetry breaking phase transitions, $SU(2) \rightarrow U(1) \rightarrow Z_2$.

The first, $SU(2) \rightarrow U(1)$, creates 't Hooft-Polyakov monopoles with mass M, the second, $U(1) \rightarrow Z_2$, confines the flux to cosmic strings with tension μ , like beads on a wire.

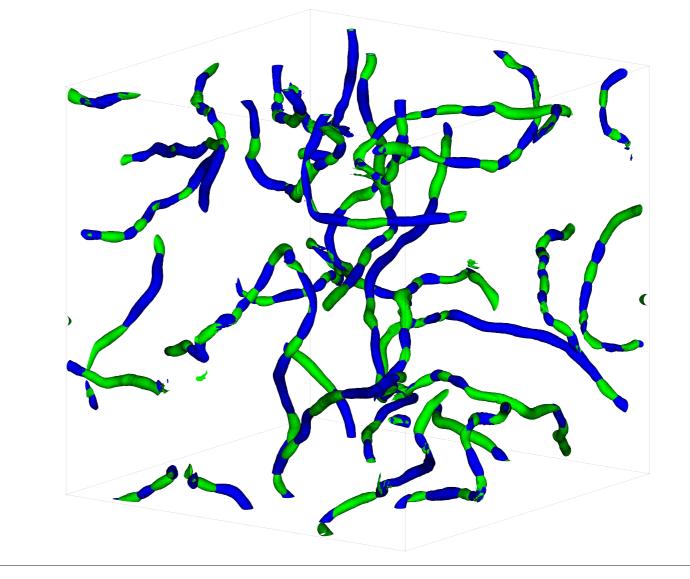
- Simulations are performed in a comoving $V = 720^3$ box with lattice spacing a = 1, with Hubble damping corresponding to an expanding radiation-dominated universe.
- We determine the location of strings and monopoles within the box, yielding L, the total (Manhattan) length of string and N, the number of monopoles. These yield the average monopole and string separations $\xi_{\rm m} = (V/N)^{1/3}$ and $\xi_{\rm s} = (V/L)^{1/2}$. The monopole locations are also used to measure \overline{v}_m , root mean square monopole velocity. • We compute monopole separation along the string d = L/N and hence the ratio $r = M/\mu d$ measuring the importance of the monopoles for the string dynamics ^{3,4,5}.

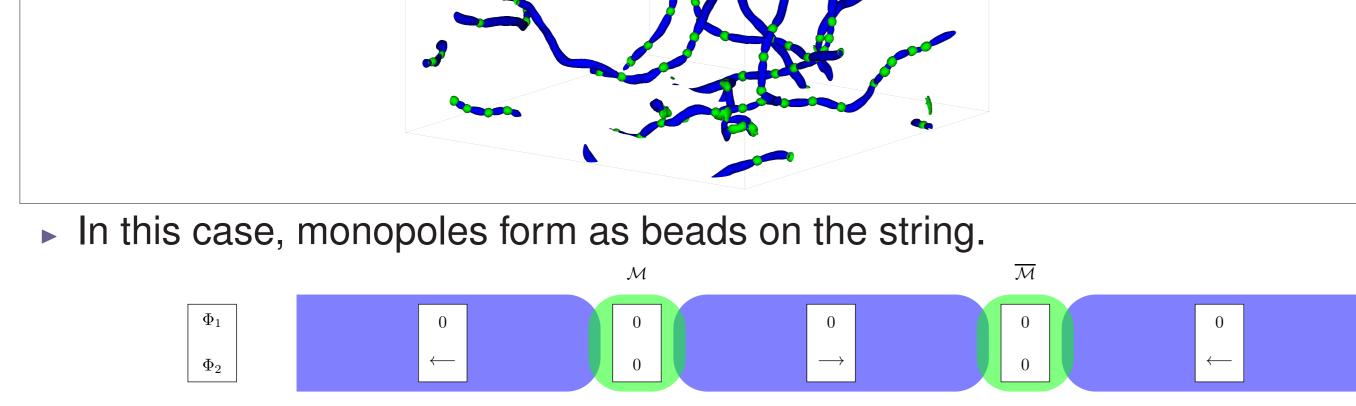
Case 1: $m_1^2 > m_2^2$

 240^3 simulation with $m_1^2 = 0.25$, $m_2^2 = 0.1$, isosurfaces $\text{Tr}\Phi_1^2 = 0.2$, $\text{Tr}\Phi_2^2 = 0.04$, time t = 240.

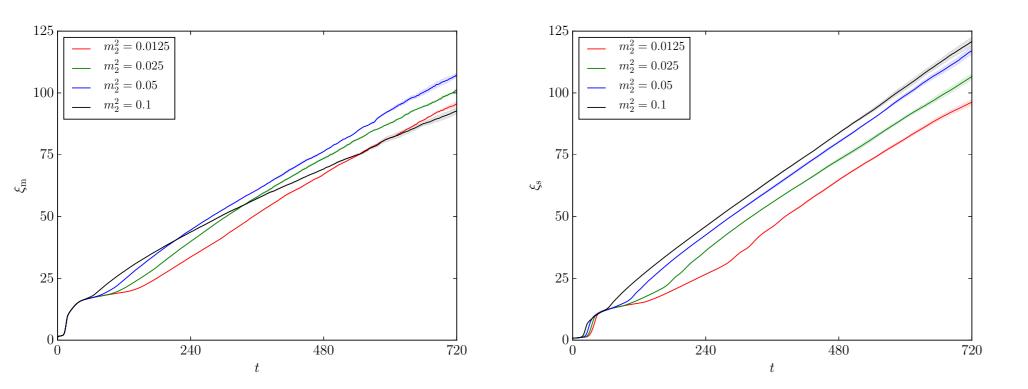
Case 2: $m_1^2 = m_2^2$

 240^3 simulation with $m_1^2 = m_2^2 = 0.25$, $\kappa = 1$, isosurfaces $\text{Tr}\Phi_1^2 = \text{Tr}\Phi_2^2 = 0.2$, time t = 240.



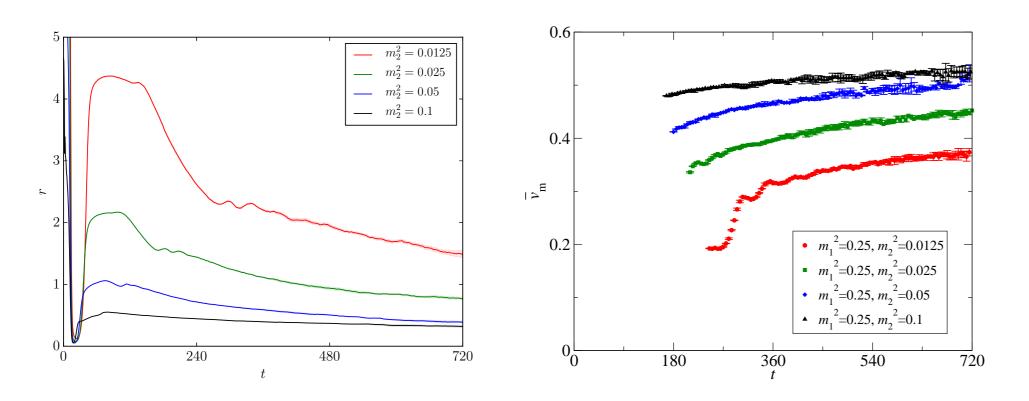


 \blacktriangleright We measured ξ_s and ξ_m and they grow linearly: a scaling network forms.

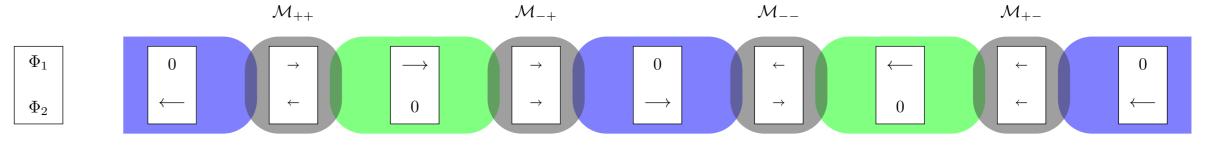


There is little apparent difference between the scaling for different μ .

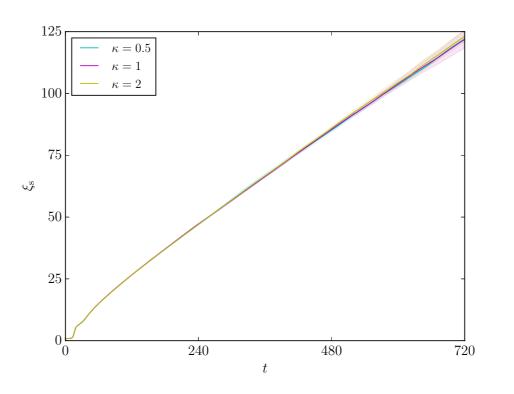
• Our simulations show that r always decreases while \overline{v}_m increases, appearing to asymptote to a relativistic value.



For all $\kappa > 0$ one gets 'half monopoles' forming at the end of string segments:



- There is a global symmetry between the two scalar fields, and the larger κ is, the less the two fields overlap.
- However, for all values of κ , the result is a scaling network of strings:



- We find that ξ_s scales with coefficient 0.16 ± 0.01 , which corresponds to string densities approximately 40% higher than in the abelian Higgs model. CMB constraints are therefore stronger for this model ⁶.
- Because of the global symmetry, we cannot count the number of half monopoles.

The monopole velocities are also in line with expected string velocities. Key results: Scaling network forms; monopoles are unimportant; average monopole velocity does not decrease.

Key results: When $m_1^2 = m_2^2$, we get novel strings with 'half monopole' structures. These still produce a scaling network.

References

1. M. Hindmarsh and T. Kibble, *Phys. Rev. Lett.* 55, 2398 (1985) 2. T. Vachaspati and A. Vilenkin, *Phys. Rev. D* 35, 1131 (1987) 3. V. Berezinsky and A. Vilenkin, *Phys. Rev. Lett.* **79**, 5202 (1997) 4. J. J. Blanco-Pillado and K. D. Olum, *JCAP* **1005**, 014 (2010) 5. C. J. A. P. Martins, *Phys. Rev. D* 82, 067301 (2010) 6. N. Bevis et al., Phys. Rev. D 82, 065004 (2010)

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